

Quantum
Mechanical &
Electromagnetic
Systems
Modelling Lab

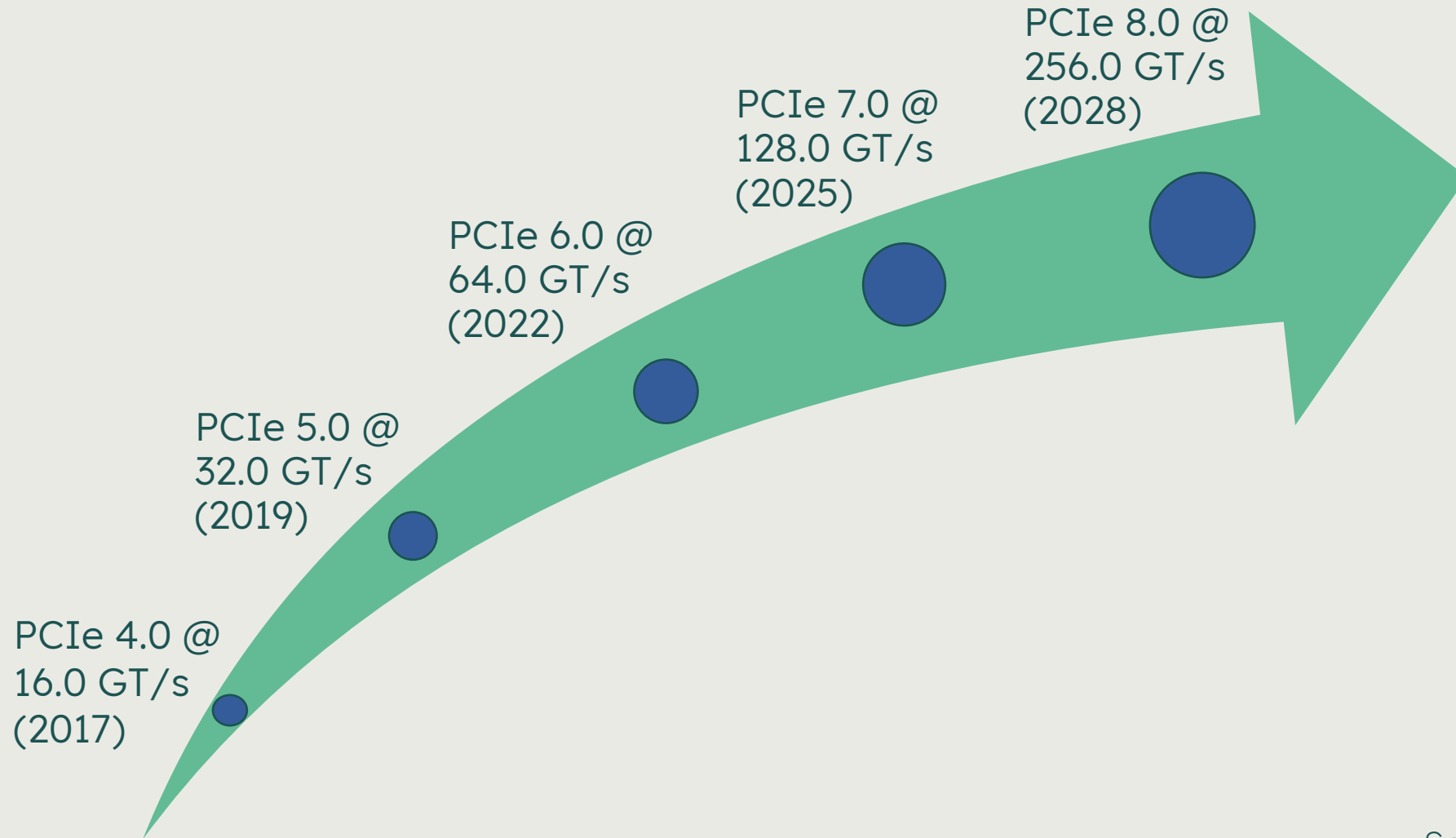
A Novel Vectorial Unified Transform for the Full-Wave Broadband Characterization of On-chip Passives

Tim Pattyn, Daniël De Zutter, Martijn Huynen, Dries Vande Ginste

quest.

Ever-increasing trend towards higher circuit speeds

Evolution of PCI Express data rate



Source: PCI-SIG,
2025



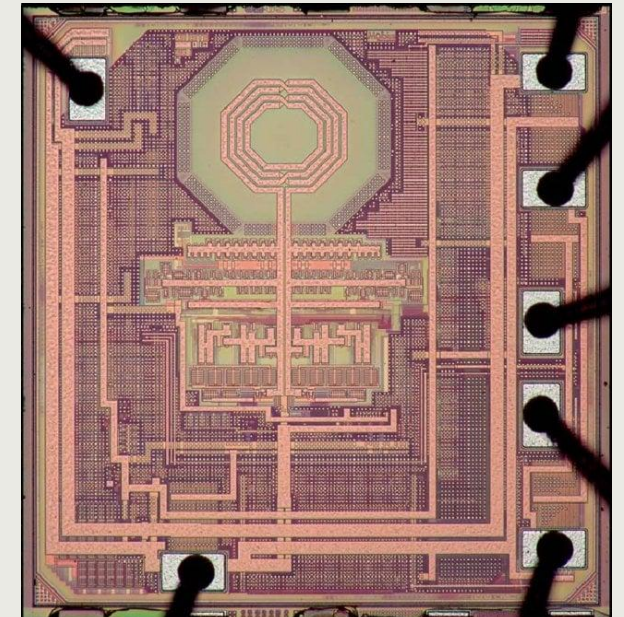
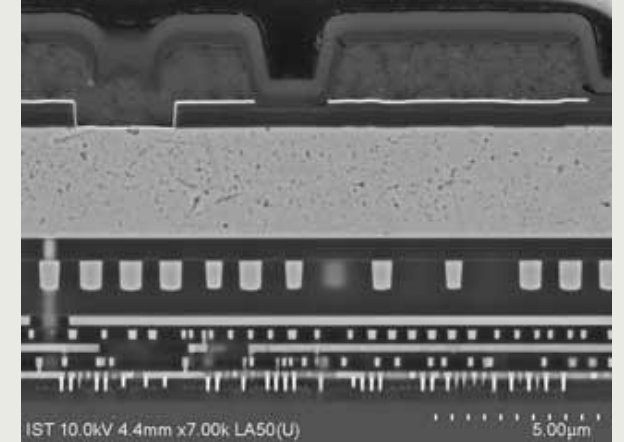
This trend requires accurate modeling

Full-wave effects become fully apparent at higher frequencies, so traditional circuit simulations no longer suffice

Heterogeneous material properties of state-of-the-art components demand precise modeling to accurately capture all electromagnetic effects

Wide variety of modeling techniques already available: FEM, VIE, PEEC, PMCHWT, GIBC, ..., but each has its own benefits and drawbacks

Differential surface admittance (DSA) formalism provides an elegant BEM framework that allows for **efficient** and **accurate** modeling, in particular broadband characterization of good conductors



At quest, we developed a novel method for constructing the 3-D DSA operator

Several techniques already available for 3-D topologies

However, existing 3-D DSA derivation methods limited to **canonical shapes** or rely on **inner region Green's function**

Leverage **unified transform framework** to alleviate both these limitations

Accurately construct DSA operator for **non-canonical shapes** without inner region Green's function



Dirichlet-to-Neumann operator using the unified transform

DSA-EFIE formulation

Numerical results

Conclusions



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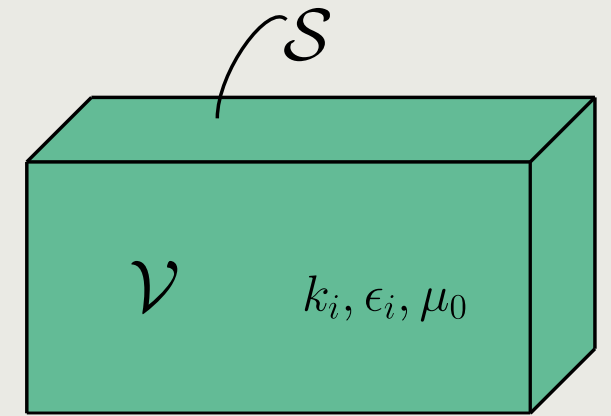
Conclusions



Dirichlet-to-Neumann operator using the unified transform

Dirichlet-to-Neumann (DtN) operator relates the tangential electric field on \mathcal{S} with the tangential magnetic field on \mathcal{S}

In this work, DtN operator is derived using the **unified transform (UT) method**



Underlying physics governed by Maxwell's equations

Governing PDE
(underlying physics)



(Discretized)
global relation



Test solutions



DtN operator

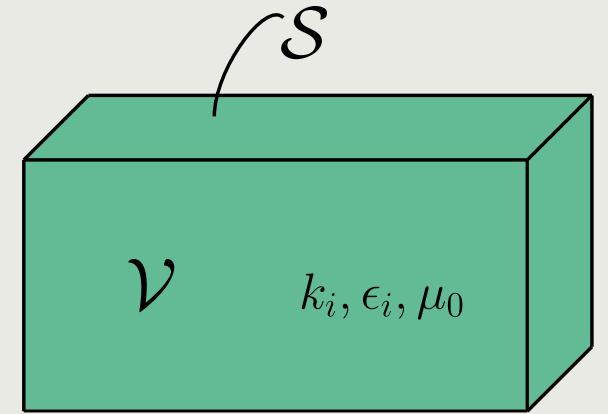
Sourceless Maxwell equations inside \mathcal{V}
Unknown solution ($\mathbf{e}_a, \mathbf{h}_a$)

$$\begin{aligned}\nabla \times \mathbf{e}_a &= -j\omega\mu_0\mathbf{h}_a \\ \nabla \times \mathbf{h}_a &= -j\omega\epsilon\mathbf{e}_a\end{aligned}$$

Known test solution ($\mathbf{e}_t, \mathbf{h}_t$)

$$\begin{aligned}\nabla \times \mathbf{e}_t &= -j\omega\mu_0\mathbf{h}_t \\ \nabla \times \mathbf{h}_t &= j\omega\epsilon\mathbf{e}_t\end{aligned}$$

Sought-after **DtN operator** relates tangential components of \mathbf{e}_a and \mathbf{h}_a on \mathcal{S}



Lorentz reciprocity as the global relation

Governing PDE
(underlying physics)



(Discretized)
global relation



Test solutions



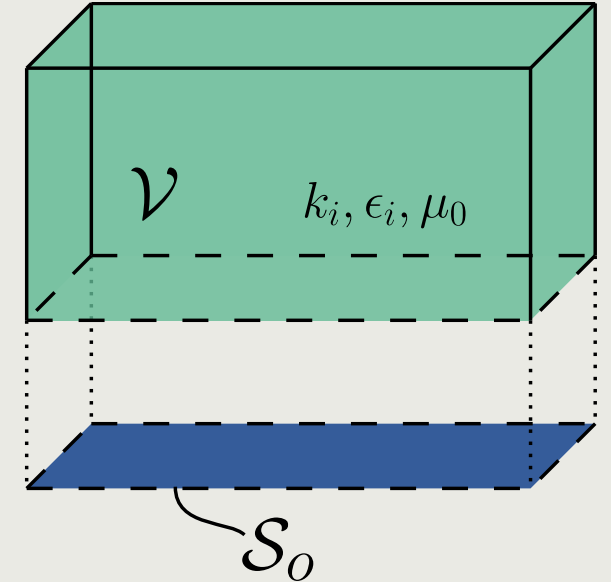
DtN operator

From **known** and **unknown** solution, we derive the global relation (= Lorentz reciprocity):

$$\iint_S (\mathbf{e}_a \cdot \hat{\mathbf{h}}_t - \mathbf{e}_t \cdot \hat{\mathbf{h}}_a) dS = 0$$

For numerical evaluation, the **unknown** solution is discretized using Legendre polynomial-based basis functions, e.g., on \mathcal{S}_0

$$\begin{aligned} & \sum_{n,m} \iint_{\mathcal{S}_0} \left((e_{nm}^{o,x} \mathbf{u}_x + e_{nm}^{o,y} \mathbf{u}_y) \cdot \hat{\mathbf{h}}_t \right) f_{nm}(x - x_o, y - y_o) dS \\ & - \sum_{n',m'} \iint_{\mathcal{S}_0} \left(\mathbf{e}_t \cdot (\hat{h}_{n'm'}^{o,x} \mathbf{u}_x + \hat{h}_{n'm'}^{o,y} \mathbf{u}_y) \right) f_{n'm'}(x - x_o, y - y_o) dS \end{aligned}$$



Test global relation using plane waves

Governing PDE
(underlying physics)



(Discretized)
global relation



Test solutions



DtN operator

Evaluate discretized global relation using two orthogonal sets of plane waves originating from \mathcal{S}_s as test solution $(\mathbf{e}_t, \mathbf{h}_t)$

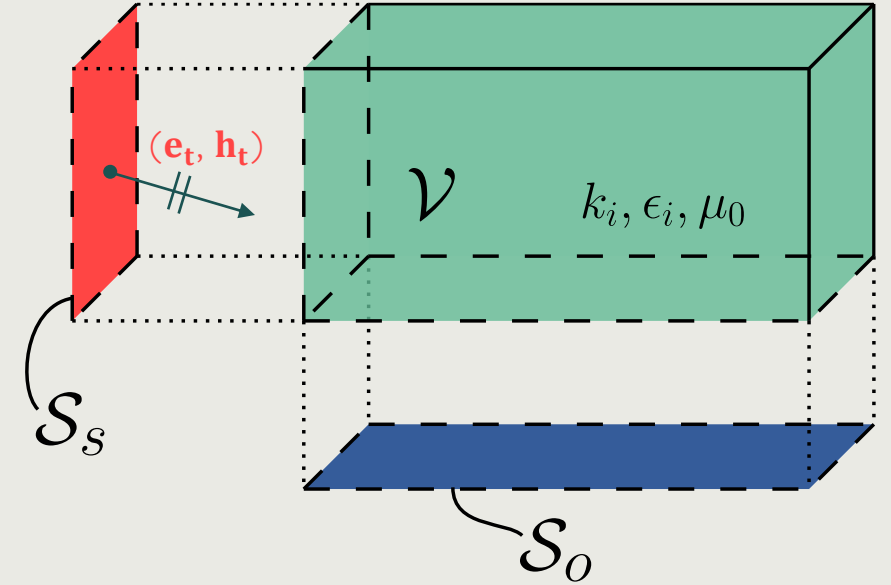
$$\text{set 1: } \mathbf{e}_t^1 = \mathbf{p}_1 e^{-j\phi}, \hat{\mathbf{h}}_1 = (\mathbf{u} \times \mathbf{p}_1) e^{-j\phi}$$

$$\text{set 2: } \mathbf{e}_t^2 = \mathbf{p}_2 e^{-j\phi}, \hat{\mathbf{h}}_2 = (\mathbf{u} \times \mathbf{p}_2) e^{-j\phi}$$

For a single wave of set 1, this results in

$$\sum_{n,m} \left(\frac{k_x k_y}{k\tau} e_{nm}^{o,x} + \frac{\tau}{k} e_{nm}^{o,y} \right) A_{nm}^{so} - \sum_{n',m'} \frac{k_z}{\tau} \hat{h}_{n'm'}^{o,y} A_{n'm'}^{so}$$

Due to Fourier-type sampling of the two sets of plane waves, a total of $2 \times (2M_{\max} + 1)(2P_{\max} + 1)$ equations is obtained



Evaluate global relation for all cuboid faces

Governing PDE
(underlying physics)



(Discretized)
global relation



Test solutions



DtN operator

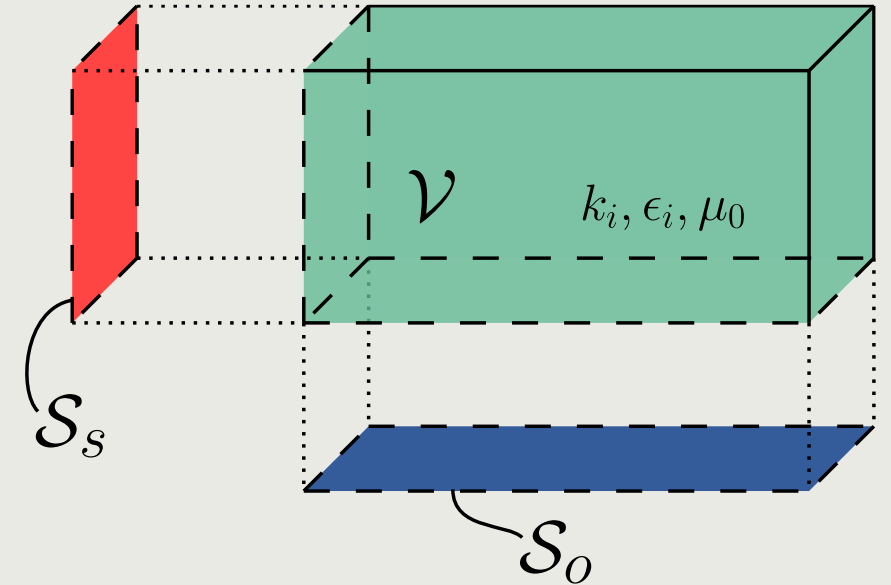
Performing this procedure for all combinations of **source** and **observer** face and collecting all equations results in

$$\bar{\mathbf{D}} \cdot \bar{\mathbf{e}} - \bar{\mathbf{N}} \cdot \bar{\mathbf{h}} = 0$$

Solving this in a least-squares sense yields the sought-after **discretized DtN operator** $\bar{\mathcal{P}}$

$$\bar{\mathbf{h}} = \bar{\mathcal{P}} \cdot \bar{\mathbf{e}}$$

No eigenmodes of \mathcal{V} or inner region Green's function required to construct $\bar{\mathcal{P}}$



Dirichlet-to-Neumann operator using the unified transform

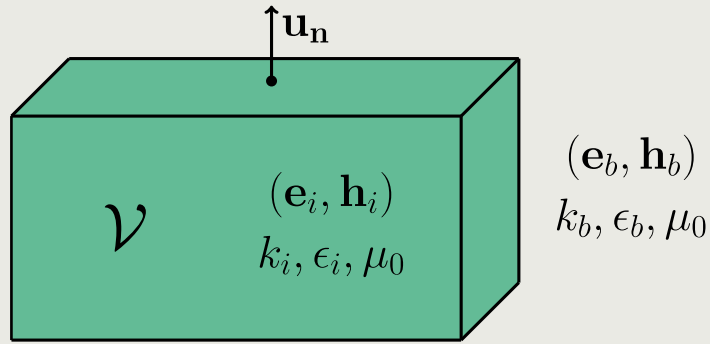
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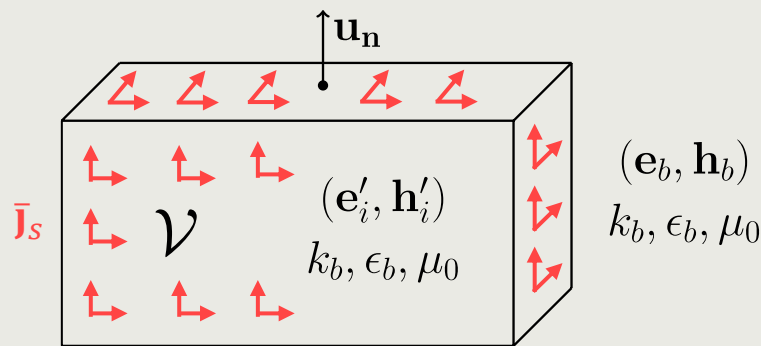
Construction of discretized DSA operator



The **equivalence principle** replaces materials by the background medium

$$\bar{\mathbf{J}}_s = \bar{\mathbf{h}}_i - \bar{\mathbf{h}}'_i = (\bar{\mathcal{P}} - \bar{\mathcal{P}}_0) \cdot \bar{\mathbf{e}}_b = \bar{\mathcal{Y}} \cdot \bar{\mathbf{e}}_b \quad \text{on } \mathcal{S}$$

DSA operator $\bar{\mathcal{Y}}$ is constructed as difference of two DtN operators $\bar{\mathcal{P}}$ and $\bar{\mathcal{P}}_0$ for the complete volume



Now, **only current sources in background medium** → easier to solve numerically

Projecting DSA operator $\bar{\mathcal{Y}}$ onto rooftop functions and integration with the **Augmented Electric Field Integral Equation [1]** provides **full characterization** of the electromagnetic problem



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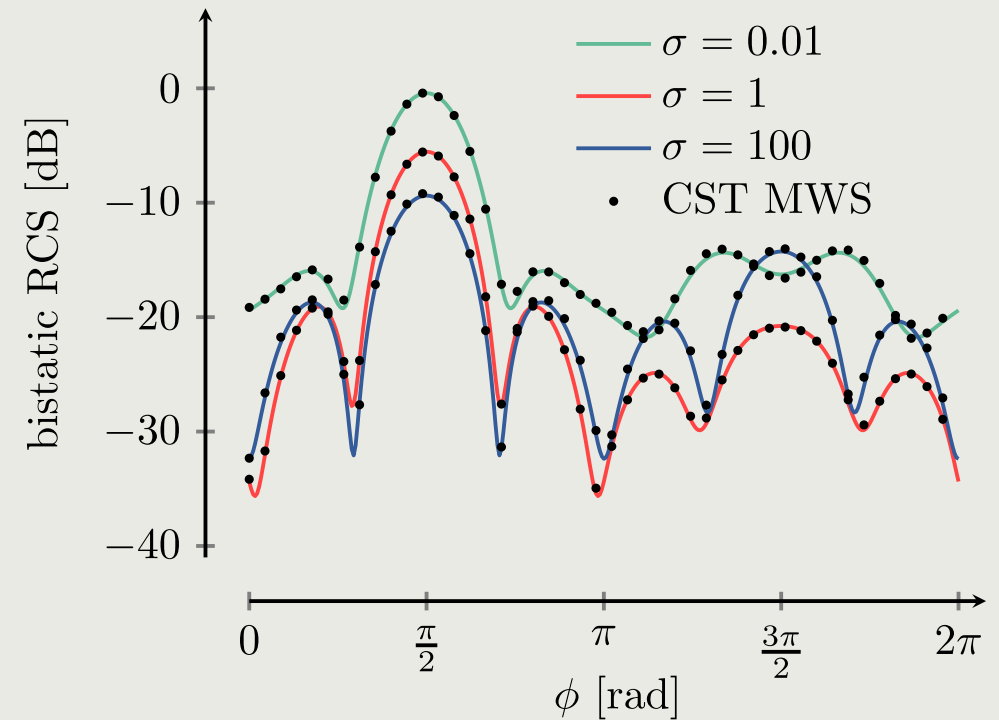
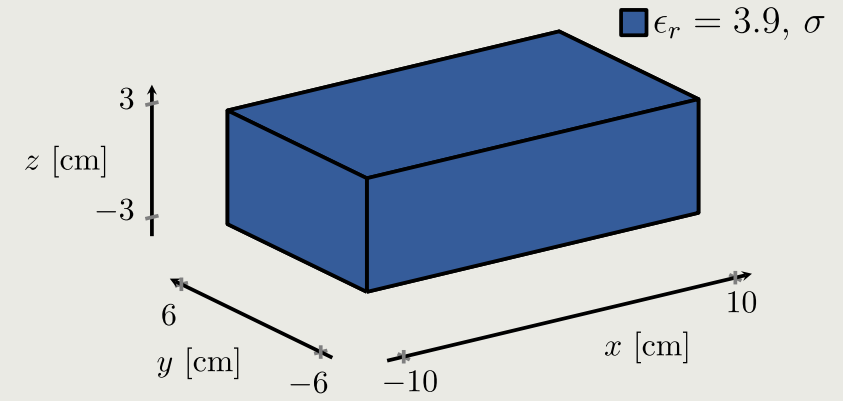
Numerical results

Dielectric cuboid with varying electrical conductivity

Illuminated by plane wave propagating in **+y-direction** with a **z-polarized** electric field

Very good agreement with CST MWS reference solution

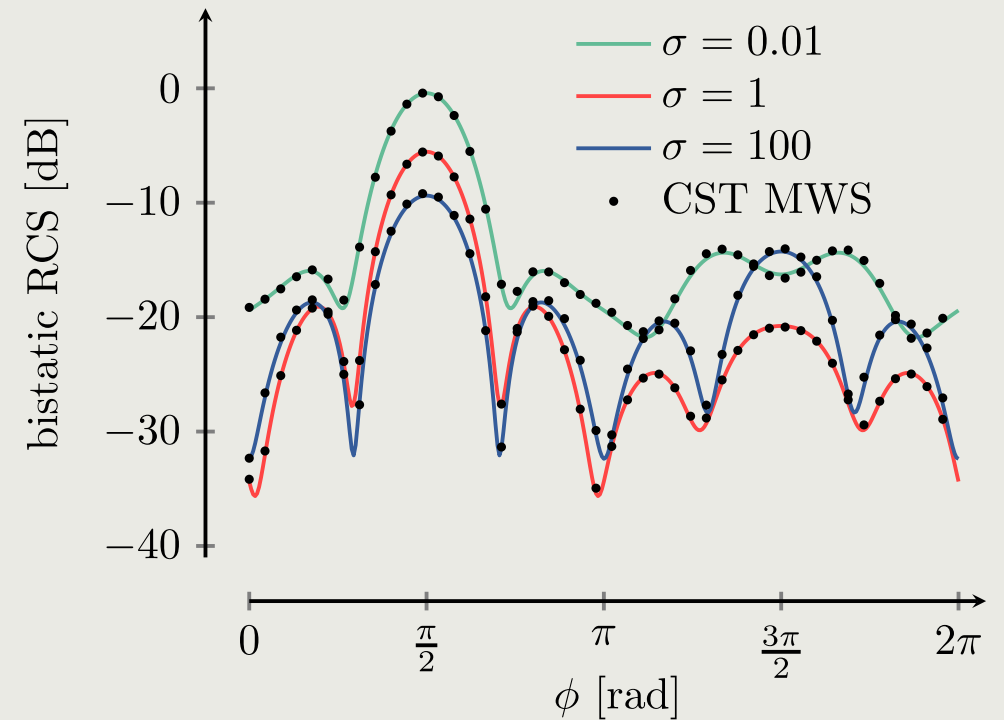
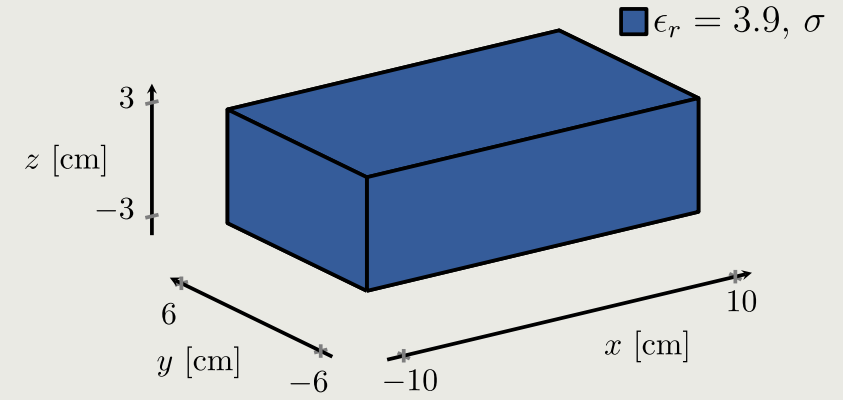
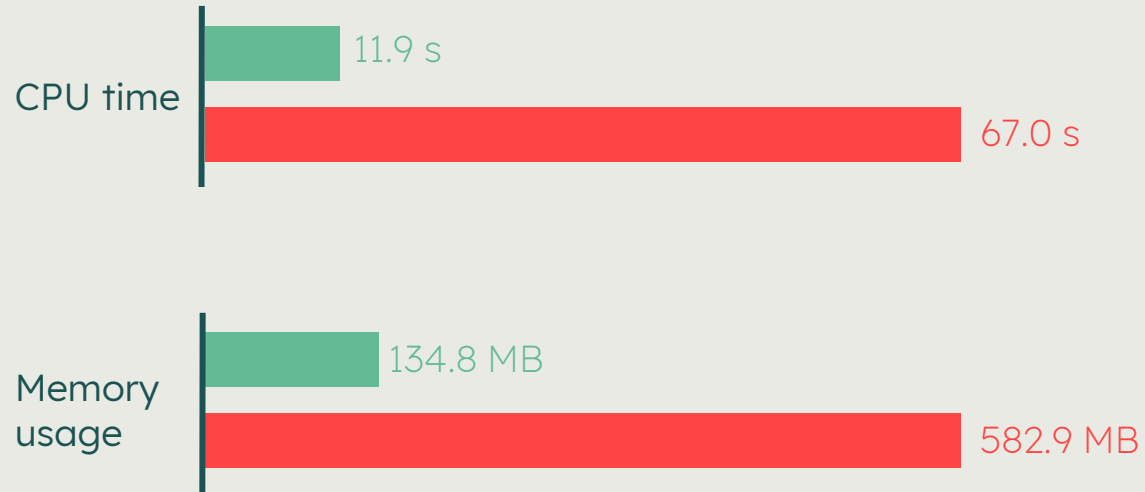
Accuracy is maintained across **wide range of material parameters**



Numerical results

Dielectric cuboid with varying electrical conductivity

Significant improvement in computational efficiency between **proposed method** and **CST MWS**



Numerical results

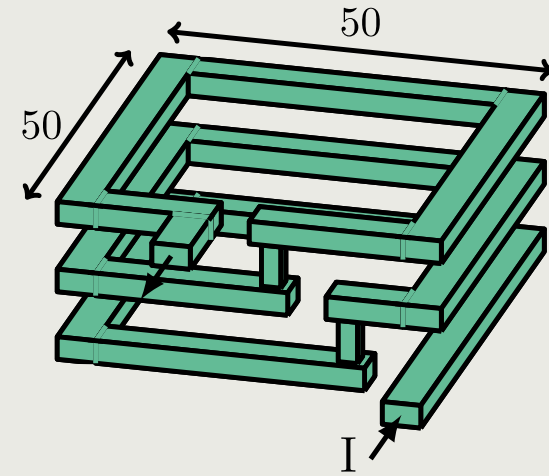
Stacked spiral on-chip inductor

Key component for various applications in radio frequency and mixed-signal integrated circuits

Material choices reflect typical on-chip integration practices

Unit current is injected into structure, allowing to determine the impedance response

- Cu ($\sigma = 5.8 \times 10^7$ S/m)
- SiO₂ ($\epsilon_r = 3.9$, $\tan\delta = 0.001$)



Numerical results

Stacked spiral on-chip inductor

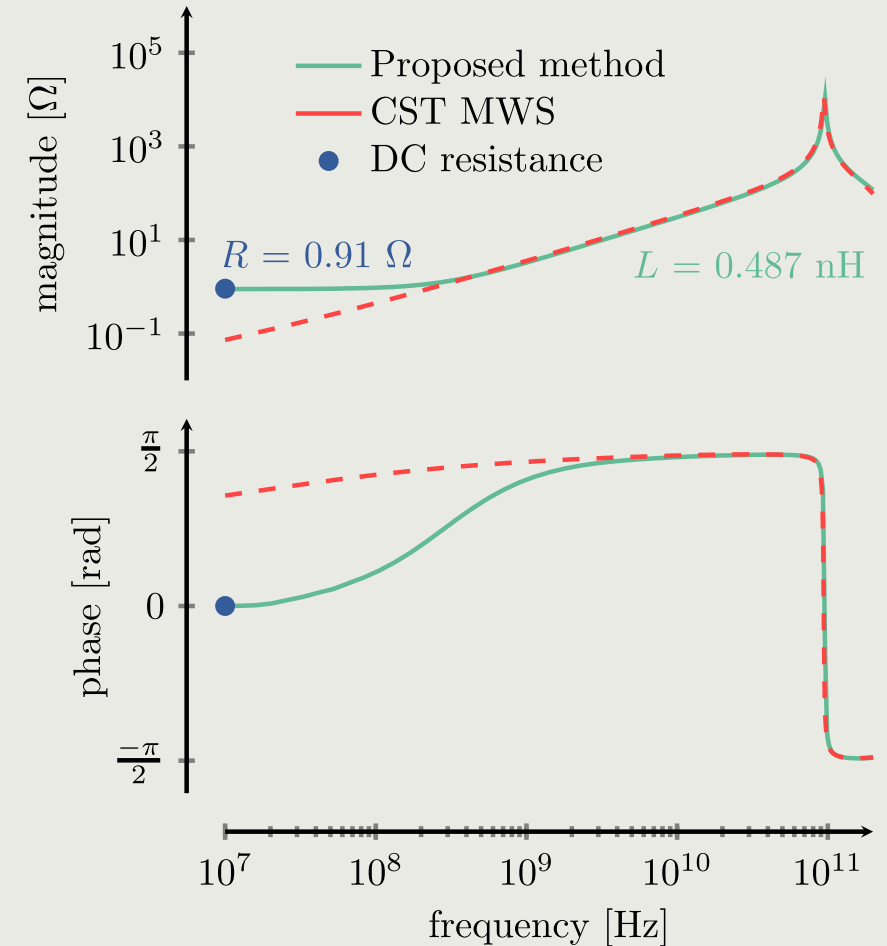
Validation with:

1. CST MWS
2. DC resistance (static solver)

Excellent agreement with CST MWS in inductive and self-resonant region

Converges to DC resistance at lower frequencies

CST MWS was unable to accurately predict the DC resistance



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Increasing circuit speed requires accurate and efficient full-wave modeling

Proposed novel unified transform method for Maxwell's equations in 3-D to determine DtN operators

DSA operator was constructed using these DtN operators and subsequently combined with the aEFIE

Demonstrated accurate broadband modeling of a wide range of materials using this new technique



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